

# Tune Spread and Tune shift induced by an electron-cloud

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# Analytical results (1)

- Equation of motion of an electron in the bunch potential (Hp. linear force)



- Evolution of the electron density during the bunch passage (→via Liouville Theorem)



- Tune shift experienced by the protons

## Analytical results (2)

- Different longitudinal bunch distribution:
  - Uniform profile
    - Gaussian distribution
  - ‘Sacherer’ distribution  $\rightarrow \sim (1-z^2)^2$
- Initially Gaussian electron distribution  $\rightarrow (\sigma_0, \sigma'_0)$

# Equation of motion of an electron in the bunch potential (1)

- Bunch distribution:

$$\rho_b(r, z) = \frac{e\lambda_b(z)}{2\pi\sigma_r^2} e^{-\frac{r^2}{2\sigma_r^2}},$$

- Field ( $\rightarrow$  from Gauss theorem)

$$E(r, z) = \frac{e\lambda_b(z)}{2\pi\epsilon_0 r} \left[ 1 - e^{-\frac{r^2}{2\sigma_r^2}} \right]$$

- Equation of motion:

$$m_e \frac{d^2 r}{dt^2} = -eE(r, z) + \frac{l^2}{m_e r^3} = -\frac{e^2 \lambda_b(z)}{2\pi\epsilon_0 r} \left[ 1 - e^{-\frac{r^2}{2\sigma_r^2}} \right] + \frac{l^2}{m_e r^3},$$

- Derivative with respect to  $\mathbf{z} = \mathbf{c} \mathbf{t}$ :

$$r'' \equiv \frac{d^2 r}{dz^2} = \frac{1}{c^2} \frac{d^2 r}{dt^2},$$

## Equation of motion of an electron in the bunch potential (2)

- x-component of the eq. of motion:

$$x'' = -\frac{2r_e \lambda_b(z)x}{r^2} \left[ 1 - e^{-\frac{r^2}{2\sigma_r^2}} \right]$$

- Approximation of linear force (for  $r \ll \sigma_r$ ):

$$x'' + \omega_e^2(z)x = 0$$

$$\omega_e^2(z) = \frac{\lambda_b(z)r_e}{\sigma_r^2}$$

- For a longitudinally uniform bunch ( $\lambda_e(\mathbf{z}) = \lambda_e$ ):

➔ harmonic oscillator:

$$\begin{aligned} x &= x_0 \cos(\omega_e z) + x'_0 \frac{1}{\omega_e} \sin(\omega_e z) \\ x' &= -x_0 \omega_e \sin(\omega_e z) + x'_0 \cos(\omega_e z) . \end{aligned}$$

# Electron density distribution

- In the **linear force approximation** the 2 planes are uncoupled  $\rightarrow$  factorization of the density in the phase space

$$\rho(x, x', y, y', z) = \rho_x(x, x', z) \rho_y(y, y', z) .$$

- Liouville Theorem + Hp. initially Gaussian Distribution

$$\rho_x(x, x', z) = \rho_x(x_0, x'_0, 0) = \frac{\sqrt{\lambda_e}}{2\pi\sigma_0\sigma'_0} e^{-\frac{x_0^2}{2\sigma_0^2}} e^{-\frac{x'^2_0}{2\sigma'^2_0}} .$$

Where  $(x_0, x'_0)$  are obtained by inverting the solution of the **eq. of motion**  $\rightarrow (\mathbf{x}_0, \mathbf{x}'_0) = \mathbf{f}(\mathbf{x}, \mathbf{x}')$

- The density is obtained by integrating:

$$n_x(x, z) = \int_{-\infty}^{+\infty} dx' \rho_x(x, x', z)$$

# Tune shift

- The field produced by the electrons is ( $\rightarrow$ Gauss theorem):

$$E(r, z) = \frac{-|e|\lambda(r, z)}{2\pi\epsilon_0 r}$$

$$\lambda(r, z) = 2\pi \int_0^r n_e(r', z) r' dr'$$

- The tune shift experienced by the protons (as a function of  $\mathbf{r}$  and  $\mathbf{z}$ ) is:

$$\Delta Q_x = \frac{1}{4\pi} \oint_C ds \beta(s) \Delta k_x$$

$$\Delta k_x = -\frac{e}{\gamma m_p c^2} \frac{\partial E_x}{\partial x}$$

# Longitudinal Uniform Profile

- eq. of motion  $\rightarrow$  harmonic oscillator:

$$\begin{aligned} x'' + \omega_e^2 x &= 0 \\ \omega_e^2 &= \frac{\bar{\lambda}_b r_e}{\sigma_r^2} \end{aligned}$$

- $(x_0, x_0') = f(x, x')$ :

$$\begin{aligned} x_0 &= x \cos(\omega_e z) - x' \frac{1}{\omega_e} \sin(\omega_e z) = Cx - \frac{S}{\omega_e} x' \\ x_0' &= x \omega_e \sin(\omega_e z) + x' \cos(\omega_e z) = \omega_e Sx + Cx' \end{aligned}$$

- electron density:

$$\begin{aligned} n(x, y, z) &= n_x(x, z) n_y(y, z) \\ &= \frac{\lambda_e \omega_e^2}{2\pi} \frac{e^{-\frac{\omega_e^2 r^2}{2(\sigma_0^2 S^2 + \omega_e^2 \sigma_0^2 C^2)}}}{\sigma_0^2 S^2 + \omega_e^2 \sigma_0^2 C^2} \end{aligned}$$

- tune shift:

$$\Delta Q_x(r, z) \approx \frac{\bar{\beta} \bar{C} \lambda_e r_p}{4\pi \gamma C^2 \sigma_0^2} \frac{1}{1 + \frac{S^2 \sigma_0^2}{C^2 \omega_e^2 \sigma_0^2}} \left( 1 - \frac{\omega_e^2 r^2}{2(C^2 \omega_e^2 \sigma_0^2 + S^2 \sigma_0^2)} \right)$$

- Hp.  $\sigma_0' \ll \sigma_0 \omega_e$ ; for  $\mathbf{r} = \mathbf{0}, \mathbf{z} = \mathbf{0}$ :

$$\Delta Q_x(r, z) \approx \frac{\bar{\beta} \bar{C} \lambda_e r_p}{4\pi \gamma \sigma_0^2}$$



# General longitudinal distribution

- eq. of motion:

$$x'' + \omega_e^2(z)x = 0$$

$$\omega_e^2(z) = \frac{\lambda_b(z)r_e}{\sigma_r^2}$$

- solution of eq. of motion in the form:

$$x(z) = A(z)e^{iS(z)}$$

- **WKB approximation:**

$$\left| \frac{3}{2} \frac{\omega_e'}{\omega_e} \right|, \left| \frac{\omega_e''}{\omega_e'} \right| \ll 2\omega_e$$



$$A(z) \approx \frac{1}{\sqrt{S'(z)}}$$

$$S'(z) \approx \omega_e(z)$$

- The general solution can also be written as:

$$x(z) = \frac{A}{\sqrt{\omega_e(z)}} \cos[S(z)] + \frac{B}{\sqrt{\omega_e(z)}} \sin[S(z)]$$

# General longitudinal distribution

- eq. of motion:

$$x'' + \omega_e^2(z)x = 0$$

$$\omega_e^2(z) = \frac{\lambda_b(z)r_e}{\sigma_r^2}$$

- $(x_0, x_0') = f(x, x')$ :

$$x_0 = a(z)x + b(z)x'$$

$$x_0' = c(z)x + d(z)x'$$

- electron density:

$$n(r, z) = \frac{\lambda_e}{2\pi[d(z)^2\sigma_0^2 + b(z)^2\sigma_0'^2]} e^{-\frac{r^2}{2[d(z)^2\sigma_0^2 + b(z)^2\sigma_0'^2]}}$$

- tune shift:

$$\Delta Q_x(r, z) \approx \frac{\bar{\beta}\bar{C}\lambda_e r_p}{4\pi\gamma [d(z)^2\sigma_0^2 + b(z)^2\sigma_0'^2]} \left( 1 - \frac{3}{4} \frac{r^2}{[d(z)^2\sigma_0^2 + b(z)^2\sigma_0'^2]} \right) + O\left(\frac{r}{\sigma_0}\right)^4$$

# 'Sacherer' distribution

- bunch longitudinal profile:

$$\lambda_b(z) = \frac{15N_b}{8l_0} \left[ 1 - \left( \frac{2z}{l_0} \right)^2 \right]^2 ; \quad z \in \left( -\frac{l_0}{2}, +\frac{l_0}{2} \right)$$

- electron density:

$$n(r, z) = \frac{\lambda_e \left( 1 - \frac{4z^2}{l_0^2} \right)}{2\pi \left\{ \cos^2 [\theta(z)] \sigma_0^2 + \frac{8\sigma_r^2 l_0}{15r_e N_b} \sin^2 [\theta(z)] \sigma_0'^2 \right\}} e^{-\frac{r^2 \left( 1 - 4z^2/l_0^2 \right)}{2 \left\{ \cos^2 [\theta(z)] \sigma_0^2 + \frac{8\sigma_r^2 l_0}{15r_e N_b} \sin^2 [\theta(z)] \sigma_0'^2 \right\}}}$$

$$\theta(z) = \sqrt{\frac{15r_e N_b}{8\sigma_r^2 l_0}} \left( z - \frac{4z^3}{3l_0^2} \right)$$

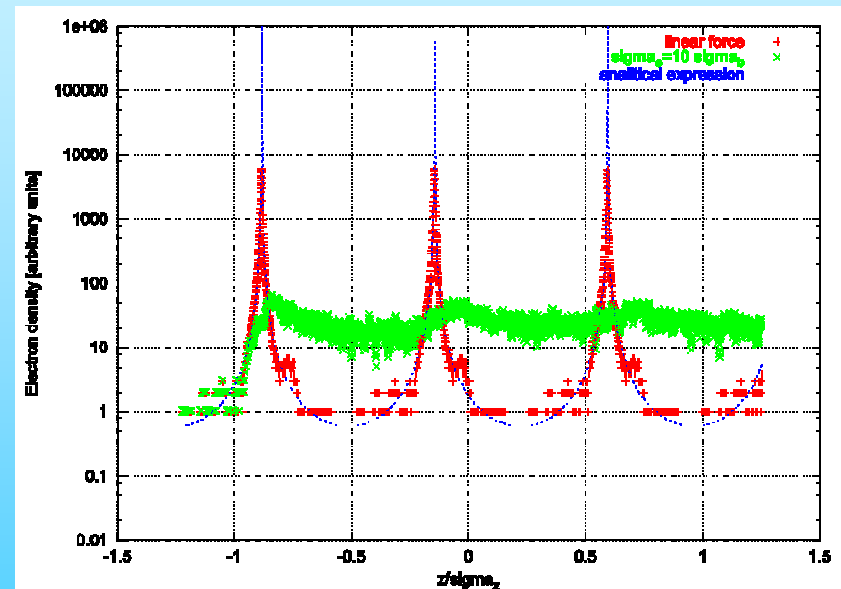
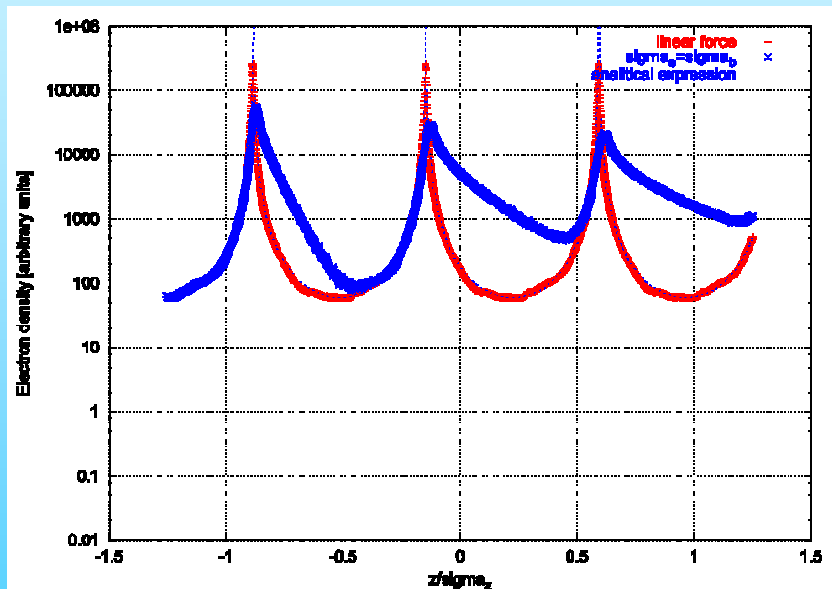
- tune shift:

$$\Delta Q_x(r, z) \approx \frac{\bar{\beta} \bar{C} \lambda_e r_p \left( 1 - \frac{4z^2}{l_0^2} \right)}{4\pi\gamma \left[ \cos^2 [\theta(z)] \sigma_0^2 - \frac{8\sigma_r^2 l_0}{15r_e N_b} \sin^2 [\theta(z)] \sigma_0'^2 \right]} \left( 1 - \frac{3}{4} \frac{\left( 1 - \frac{4z^2}{l_0^2} \right) r^2}{\left( \cos^2 [\theta(z)] \sigma_0^2 - \frac{8\sigma_r^2 l_0}{15r_e N_b} \sin^2 [\theta(z)] \sigma_0'^2 \right)} \right)$$

- Hp.  $\sigma_0' \ll \sigma_0 \omega_e$  ; for  $\mathbf{r} = \mathbf{0}$  ,  $\mathbf{z} = \mathbf{0}$ :

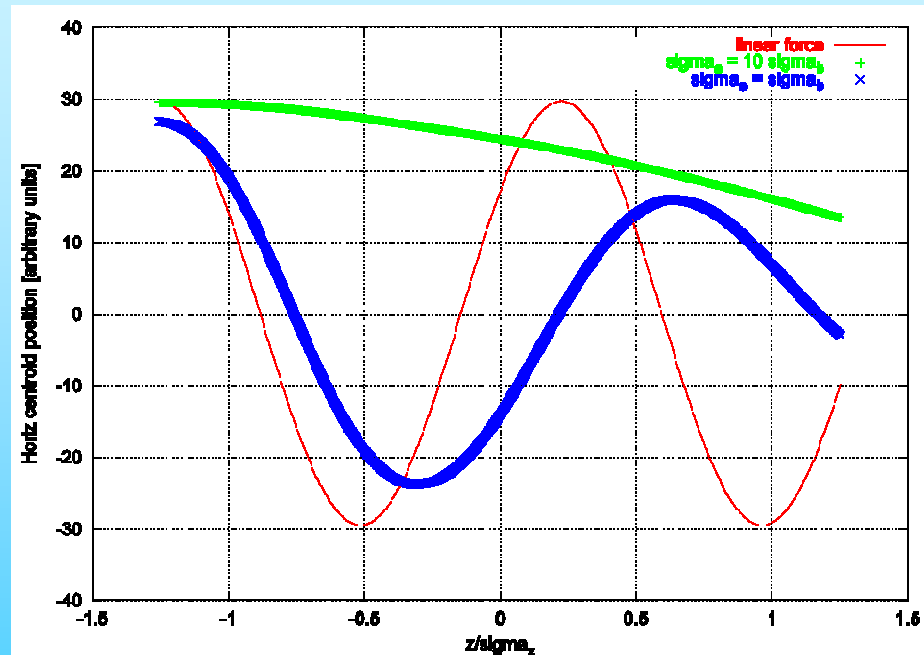
$$\Delta Q_x(r, z) \approx \frac{\bar{\beta} \bar{C} \lambda_e r_p}{4\pi\gamma \sigma_0^2}$$

# Simulations: Gaussian force + longitudinally uniform bunch



Electron density vs. time at the centre  $r=0$  for an electron size =  $\sigma_b$  ( Left) and  $10 \cdot \sigma_b$  (Right)

# Simulations: Gaussian force + longitudinally uniform bunch



Horizontal centroid position Vs. time for an electron size= $10 \sigma_b$  (green) and  $\sigma_b$  (blue), for an initial offset of  $1/\sigma_b$