

advanced theory of collimator wake
and application to SPS collimator
experiment

nonlinear resistive wall wake field

beam size effect on coherent tune shift

incoherent tune spread

Potential for nonlinear resistive-wall impedance between two parallel plates was derived by Piwinski (DESY 94-068, Eq. (52)) and re-written by Bane, Irwin, and Raubenheimer (NLC ZDR p. 594).

$$V(x, y, x_0, y_0) = -\kappa f_R(\tau) \left[\frac{-x_- \sinh x_- + y_+ \sin y_+}{\cosh x_- + \cos y_+} + \frac{x_- \sinh x_- + y_- \sin y_-}{\cosh x_- - \cos y_-} \right]$$

$$y_+ = \frac{\pi}{2g}(y + y_0), \quad y_- = \frac{\pi}{2g}(y - y_0), \quad x_- = \frac{\pi}{2g}(x - x_0) \quad 2g: \text{ full gap}$$

$$\kappa = \frac{1}{2} \frac{Nr_p}{\gamma\sigma_z} \frac{L}{g} \sqrt{\lambda\sigma_z} \quad f_R(\tau) = \frac{\sqrt{2}}{\pi} \int_0^\infty \frac{d\tau'}{\sqrt{\tau'}} e^{-\frac{(\tau+\tau')^2}{2}} \quad \langle f_R \rangle = 0.816$$

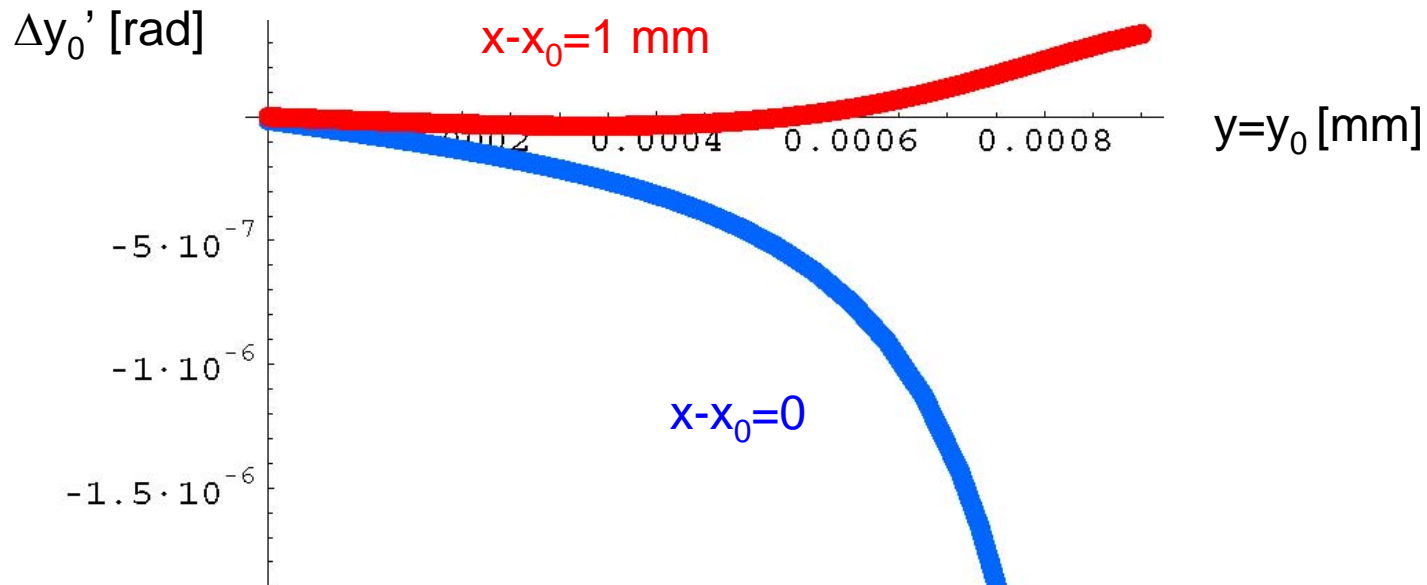
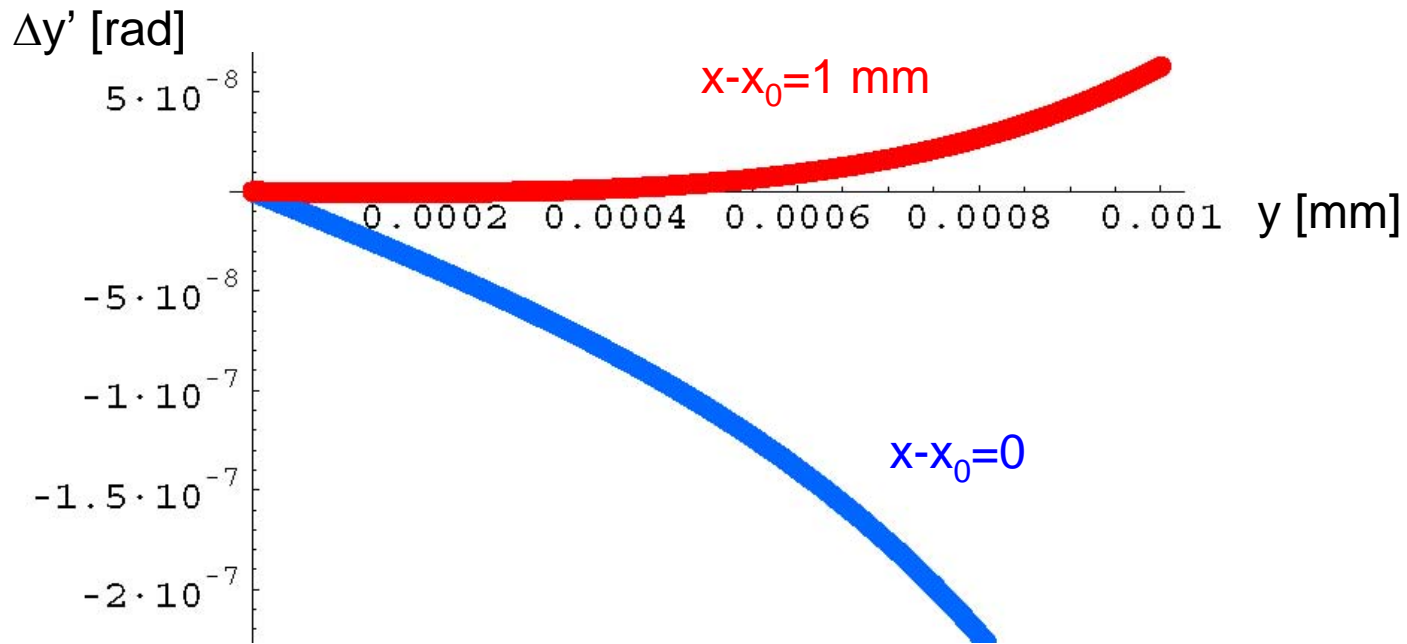
$$\lambda = \rho / (120\pi)$$

nonlinear kick to test particle:

$$\Delta y' = - \left\langle \frac{\partial V}{\partial y} \right\rangle$$

Piwinski formula applies if $\delta_s \ll \text{Min}(d_{wall}, d_{bw}, 1/(bk^2))$

many thanks to Karl Bane who helped understanding ZDR typos & incorrect reference



I calculate the coherent tune shift as follows

$$\Delta Q_{coh} = -\frac{\beta}{4\pi} \kappa \langle f_R \rangle \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-g}^g \frac{\partial(\Delta y'(\tilde{x}, y = y_0 + \tilde{y}, x_0, y_0))}{\partial y_0} \Big|_{y_0=0} \frac{e^{-\frac{\tilde{y}^2}{2\sigma_y^2}} e^{-\frac{x^2}{2\sigma_x^2}} e^{-\frac{x_0^2}{2\sigma_x^2}}}{(2\pi)^{3/2} \sigma_x^2 \sigma_y} d\tilde{y} dx dx_0$$

note: for the SPS experiment we need to exchange x and y planes

and the incoherent tune shift as

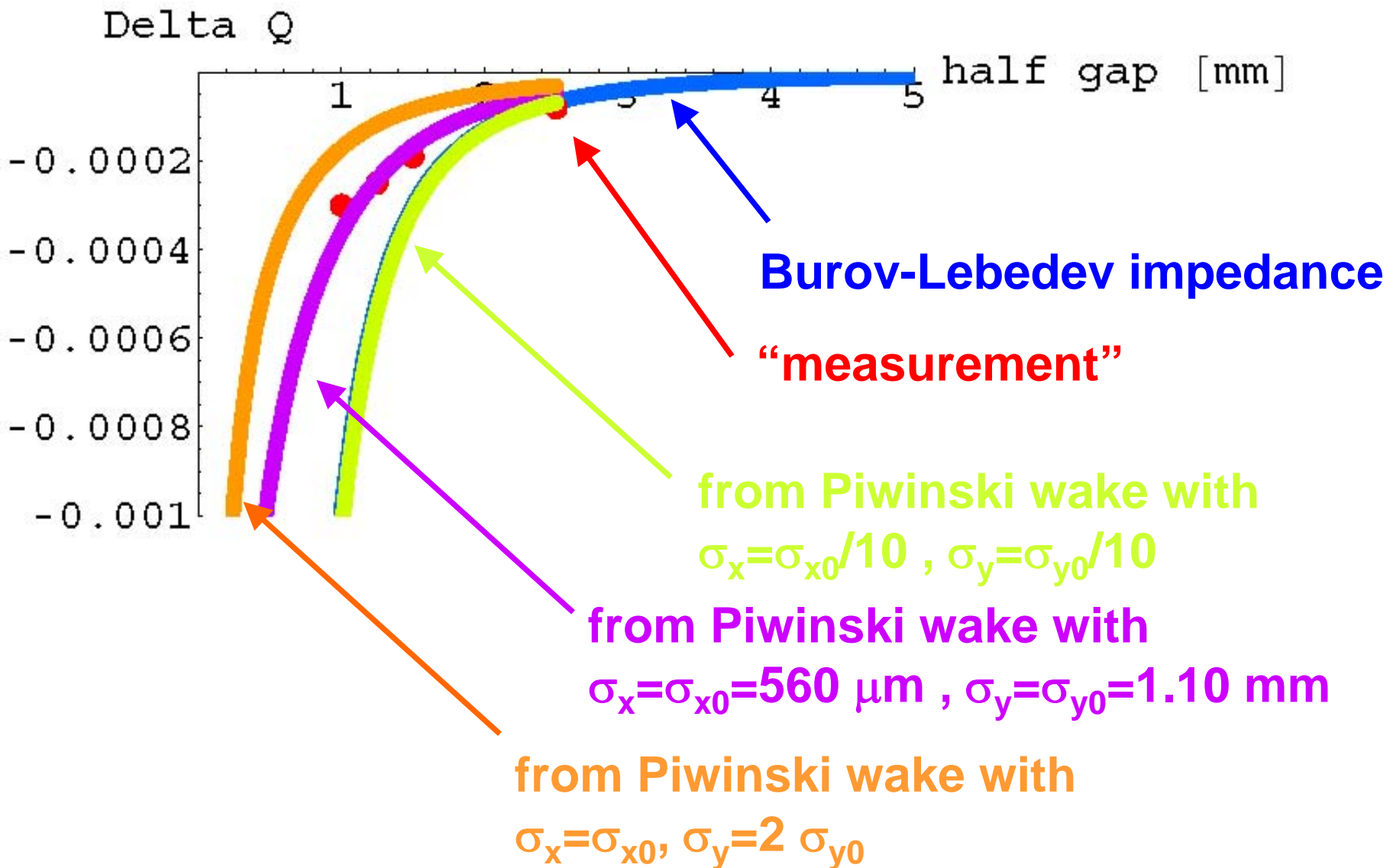
$$\langle \Delta Q_{y,inc} \rangle = -\frac{\beta}{4\pi} \kappa \langle f_R \rangle \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} \int_{-g}^g \frac{\partial(\Delta y'(\tilde{x}, y = \tilde{y}, x_0, y_0))}{\partial \tilde{y}} \Big|_{\tilde{y}=0} \frac{e^{-\frac{y_0^2}{2\sigma_y^2}} e^{-\frac{x_0^2}{2\sigma_x^2}}}{2\pi \sigma_x \sigma_y} dy_0 dx_0 \right] \frac{e^{-\frac{x^2}{2\sigma_x^2}}}{(2\pi)^{1/2} \sigma_x} dx$$

$$\langle (\Delta Q_{y,inc})^2 \rangle = \int_{-\infty}^{\infty} \left[\frac{\beta}{4\pi} \kappa \langle f_R \rangle \int_{-\infty}^{\infty} \int_{-g}^g \frac{\partial(\Delta y'(\tilde{x}, y = \tilde{y}, x_0, y_0))}{\partial \tilde{y}} \Big|_{\tilde{y}=0} \frac{e^{-\frac{y_0^2}{2\sigma_y^2}} e^{-\frac{x_0^2}{2\sigma_x^2}}}{2\pi \sigma_x \sigma_y} dy_0 dx_0 \right]^2 \frac{e^{-\frac{x^2}{2\sigma_x^2}}}{(2\pi)^{1/2} \sigma_x} dx$$

$$\Delta Q_{y,rms}^{rms} = \sqrt{\langle (\Delta Q_{y,inc})^2 \rangle - \langle \Delta Q_{y,inc} \rangle^2}$$

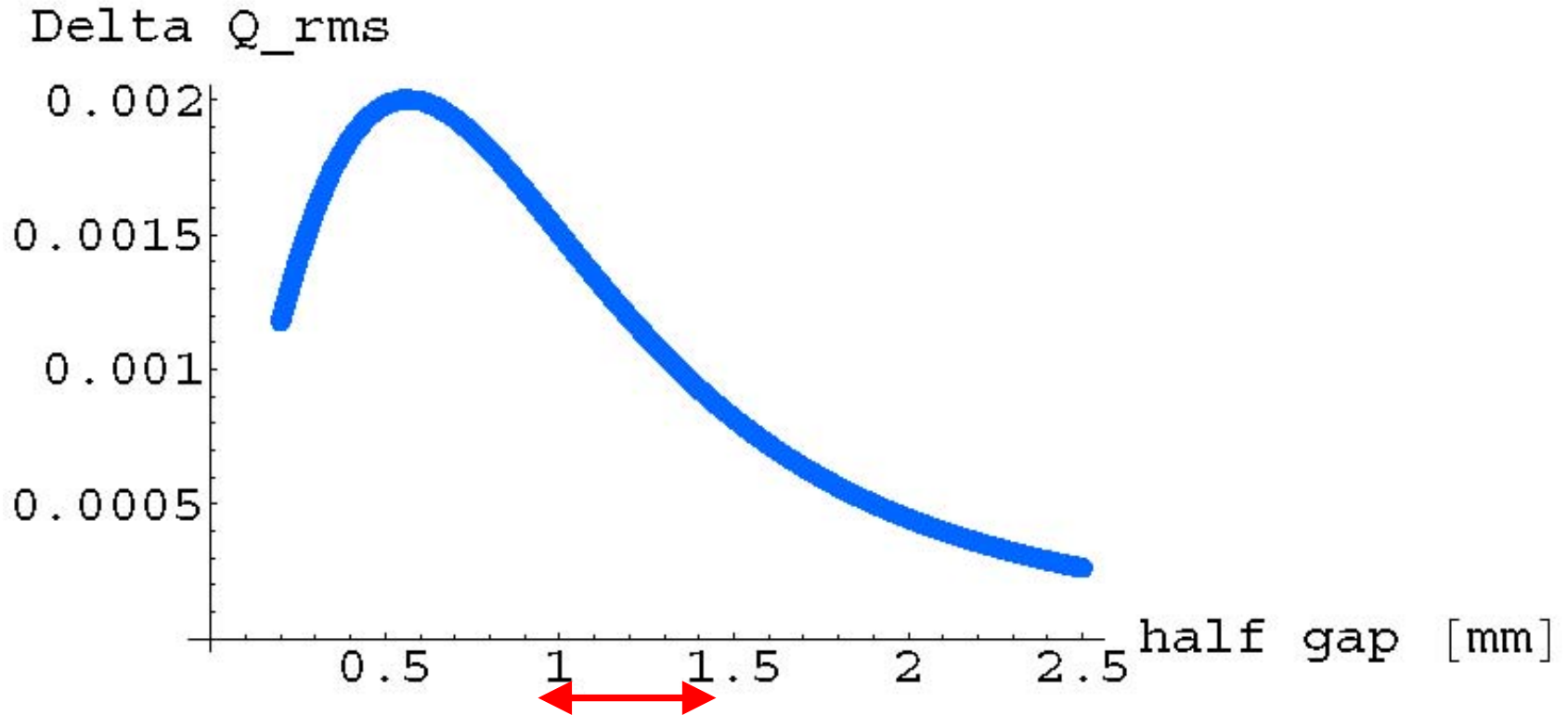
equations take scraping losses into account

coherent tune shift



the data are ~best fitted for the actual beam size estimate

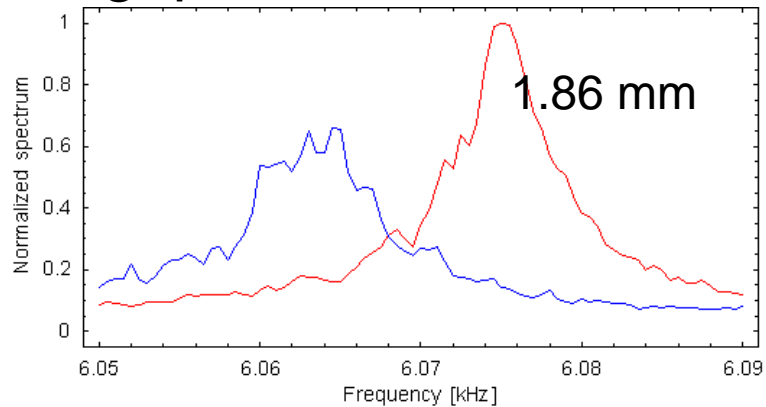
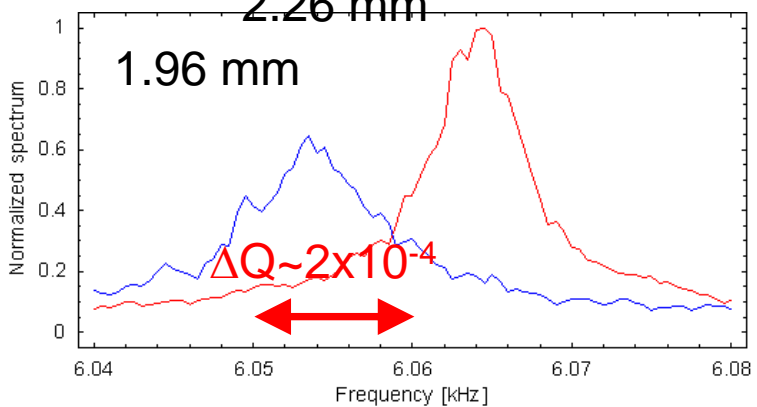
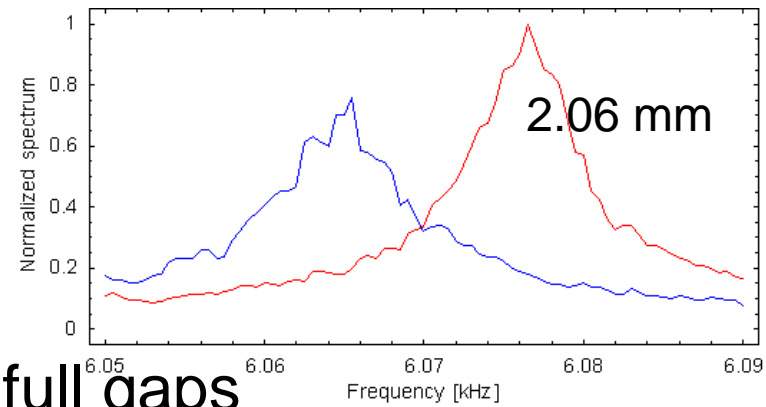
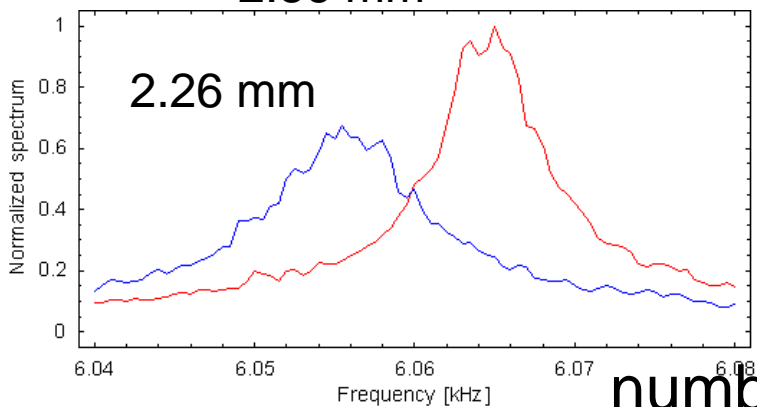
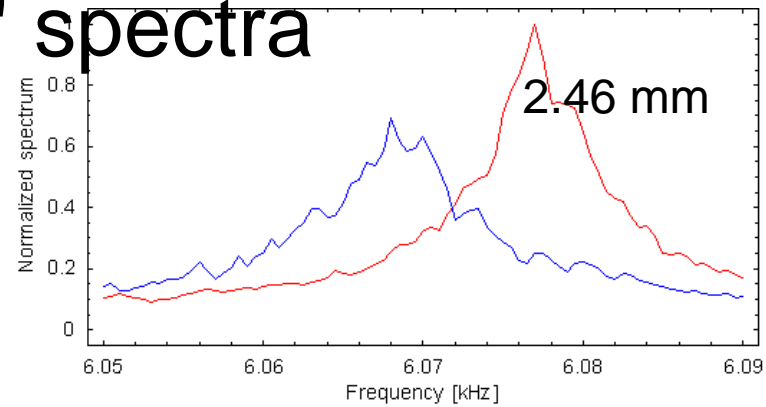
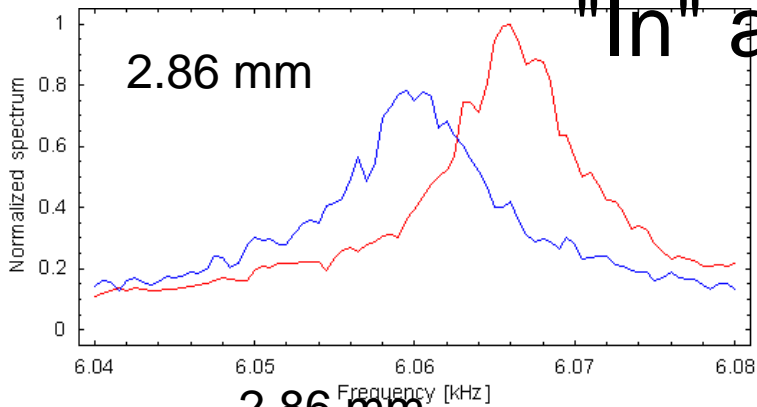
rms incoherent tune spread



(only from transverse rms, i.e., does not include longitudinal variation of f_R)

maximum value if gap \sim rms beam size

"In" and "out" spectra



number refer to full gaps