FOLLOW-UP OF THE DATA ANALYSIS FOR THE FAST VERTICAL SINGLE-BUNCH INSTABILITY AT SPS INJECTION

E. Métral

2 MDs done on

- Friday 12/09/2003
- Monday 10/11/2003
- Until now we looked essentially at the intensity threshold vs. chromaticity, etc... (analytical estimate + simulations with MOSES and HEADTAIL codes)
- Now we want to analyze in detail the time evolution (some measurements have be made with the HEADTAIL monitor)

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MEASUREMENTS FROM THE HEADTAIL MONITOR (1/6)

\Rightarrow Data taken and 1st treated by G. Arduini (sent to me on 10/08/05)

Example here (12/09/03): ~ 1.2 10¹¹ p, ~ 0.2 eVs, ξ_ν ~ 0.05, T_s ~ 300 turns



MEASUREMENTS FROM THE HEADTAIL MONITOR (2/6)



MEASUREMENTS FROM THE HEADTAIL MONITOR (3/6)

FFT applied over the full acquisition depth (372 turns ~ 8.6 ms). The peaks at 40 MHz are due to the fact that at each turn only 25 ns out of the full revolution period of 23.1 μ s are acquired



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MEASUREMENTS FROM THE HEADTAIL MONITOR (4/6)

Betatron phase difference between different temporal slices with respect to the central slice (=centre of the bunch)



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MEASUREMENTS FROM THE HEADTAIL MONITOR (5/6)



MEASUREMENTS FROM THE HEADTAIL MONITOR (6/6)

Tune deduced at each turn from data analysis (using Hilbert transform techniques)



HEAD-TAIL PHASE SHIFT THEORY (1/5)

Vertical excursion of a slice of a bunch after a vertical kick (linear model) \implies Cf. LHC Project Report 602 (Fartoukh&Jones)

$$\langle y \rangle (\tau, n) = A(\tau, n) \sin \left[2 \pi n Q_y + \phi_\beta(\tau, n) \right]$$

Number of machine turns

Time delay w.r.t central slice $\tau = 0$

with

$$A(\tau, n) = A e^{-\frac{\left(\frac{\omega_0 Q'_y}{\eta} \sigma_\tau \sin\left[2\pi n Q_s\right]\right)^2}{2}}$$

$$\phi_{\beta}(\tau,n) = \frac{\omega_0 Q'_y}{\eta} \tau \left(1 - \cos\left[2\pi n Q_s\right]\right)$$

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HEAD-TAIL PHASE SHIFT THEORY (3/5)



HEAD-TAIL PHASE SHIFT THEORY (4/5)



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HEAD-TAIL PHASE SHIFT THEORY (5/5)



HEAD-TAIL & TMC INSTABILITY THEORY (1/27) Computation with MOSES for 0 chromaticity

- Real Part of $(v-v_x)/v_s$ -MOSES -- MODE COUPLING INSTABILITY IN SPS AT 26 GEV 23/08/05 14.26.09 VERSION 3.3 CPU TIME USED: 0.531-314 (s) 0 SPRD = 0.000E+00NUS = 0.324E-02 ENGY = 26.0 (GeV) SGMZ = 21.0 (cm) $BETAC = 40.0 \quad (m)$ REVFRQ= 0.433E-01 (MHz) -2 ALPHA = 0.192E-02 CHORM = 0.000E+00 FREQ = 700. (MHz) RS = 10.0 (MOhm/m) OV = 1.00 LBIN = F-4 Real $(v-v_X)/v_S$ MU = 5 **BB** resonator impedance -6 $R_v = 10 \text{ M}\Omega/\text{m}$ $f_r = 700 \text{ MHz}$ -8 0.5 0.25 0.75 0 1 $I_{h}(mA)$ O = 1

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HEAD-TAIL & TMC INSTABILITY THEORY (2/27)



HEAD-TAIL & TMC INSTABILITY THEORY (3/27)

If we suppose an initial offset of ~ 1mm and assume that beam losses occur after a growth of a factor ~ 20 (same approximation as for the BBU mechanism discussed later), then beam losses should occur after ~ 3 rise-times, i.e. after ~ 130 turns

HEAD-TAIL & TMC INSTABILITY THEORY (4/27) Reminder: Simple formula to have a rough estimate of the TMC instability rise-time (for 0 chromaticity)

$$\tau_{TMC} = \frac{T_s}{\pi \sqrt{\left(\frac{N_b}{N_b^{th}} - 1\right)\left(\frac{N_b}{N_b^{th}}q + 1\right)}}$$
with $q \in [0,1]$ $q = 0$ for short bunch, i.e. $2 f_r \tau_b \approx 1$
 $q = 1$ for long bunch, i.e. $2 f_r \tau_b \gg 1$
Numerical application $\tau_{TMC} [q = 0] = 66$ SPS turns
 $\tau_{TMC} [q = 1] = 32$ SPS turns
 $\tau_{TMC} [q = 0.5] = 41$ SPS turns
To MOSES

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HEAD-TAIL & TMC INSTABILITY THEORY (5/27)



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HEAD-TAIL & TMC INSTABILITY THEORY (6/27)



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HEAD-TAIL & TMC INSTABILITY THEORY (7/27)

Using the simple formula one has

$$2 f_r \tau_b = 7.3$$

$$\tau_{TMC} [q=0] = 158$$
 SPS turns
 $\tau_{TMC} [q=1] = 102$ SPS turns
 $\tau_{TMC} [q=0.5] = 121$ SPS turns
 $\tau_{TMC} [q=0.5] = 121$ SPS turns

HEAD-TAIL & TMC INSTABILITY THEORY (8/27) Computation with MOSES for nonzero chromaticities



HEAD-TAIL & TMC INSTABILITY THEORY (9/27)



HEAD-TAIL & TMC INSTABILITY THEORY (10/27)



HEAD-TAIL & TMC INSTABILITY THEORY (11/27)



HEAD-TAIL & TMC INSTABILITY THEORY (12/27)



HEAD-TAIL & TMC INSTABILITY THEORY (13/27)



HEAD-TAIL & TMC INSTABILITY THEORY (14/27)



HEAD-TAIL & TMC INSTABILITY THEORY (15/27)



HEAD-TAIL & TMC INSTABILITY THEORY (16/27)



HEAD-TAIL & TMC INSTABILITY THEORY (17/27)



HEAD-TAIL & TMC INSTABILITY THEORY (18/27)





HEAD-TAIL & TMC INSTABILITY THEORY (20/27)

IF the different head-tail modes could be treated separately (i.e. below the TMC threshold), one would expect vertical signals as below (assuming the same rise-time from MOSES)





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HEAD-TAIL & TMC INSTABILITY THEORY (23/27)

• For mode

m





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HEAD-TAIL & TMC INSTABILITY THEORY (25/27)

For mode

m





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HEAD-TAIL & TMC INSTABILITY THEORY (27/27)



BBU INSTABILITY THEORY (1/9)

Yokoya DESY 86-084, Brandt&Gareyte (CERN SPS/88-17 (AMS))



BBU INSTABILITY THEORY (2/9)



BBU INSTABILITY THEORY (3/9)



BBU INSTABILITY THEORY (4/9)



BBU INSTABILITY THEORY (5/9)





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BBU INSTABILITY THEORY (6/9)



BBU INSTABILITY THEORY (7/9)



BBU INSTABILITY THEORY (8/9)



BBU INSTABILITY THEORY (9/9)



CONCLUSION (1/5)

 The time evolution from injection to beam losses seems to be in good agreement with both TMC and BBU theory (for low chromaticity)

Assuming a Broad-Band resonator impedance with

$$R_y = 10 \text{ M}\Omega/\text{m}$$
 $f_r = 700 \text{ MHz}$ $Q = 1$

yields beam losses after

- ~ 130 SPS turns from the TMC theory
- ~ 105 SPS turns from the BBU theory

whereas during the MD beam losses occurred after ~ 100-150 turns

CONCLUSION (2/5)

The BBU theory was already used in 1988 by Gareyte&Brandt to explain the fast (11 turns) instability of positron bunches observed at SPS injection



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CONCLUSION (3/5)

More detailed simulations are planned with the HEADTAIL code. One case already studied (with similar beam parameters) is shown below (cf. CERN-AB-2004-055). This picture is similar to the one obtained with the BBU formalism (careful check of parameters!...)



CONCLUSION (4/5)

- The TMC instability formalism clearly shows the increase of the instability rise-time with chromaticity, as observed qualitatively.
 A quantitative comparison remains to be done
- The analysis of the betatron phase differences between different temporal slices with respect to the central slice (=centre of the bunch) revealed interesting features, as seen below



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